An AXIOMATIC ON PROXIMITIES António G. PORTELA CAUL, Complexo II INIC, Av. Prof. Gama Pinto, 2 1699 BISBOA Codex Summary: An introduction of the concept of Proximity, namely C-proximities, is given, a typical axiomatic is presented, and an application to thermodynamics shows that C-proximities are rather flexible. This is clearly a preliminar presentation but we hope to have shown the power of fuzzy sets in connections to C-prox. Keywords: proximity, fuzzy Sets, unimodal membership function, thermodynamics, homogeneity, entropy measures.

D - INTRODUCTION

PROXIMITY (PROX) was introduced (11), to deal with real problems that di not "fit" in metrizable spaces, and even in pseudo-metrics of different kinds were not suitable.

Essentially PROX is a function \mathcal{H} (eventually a generalized one) defined on X × X and taking values in \mathcal{G} , when X is the "formal set" where the images of the real problem are projected and \mathcal{G} is a suitable set.

 \mathcal{H} is defined in such a way that a <u>loose</u> triangular law is retained. Here, the elements of \mathcal{G} are fuzzy sets, with characterizers Γ which are of a special kind, namely <u>unimodal</u> the justification for this name is the shape of Γ that is similar to the unimodal distributions.

Proximities based on the axiomatic introduced in Chapter 1 and using unimodal characterizers were named $\,\mathcal{C}$ -Proximities, $\,\mathcal{C}$ -PROX for short. A topology on $\,\chi\,$ is induced through $\,\mathcal{H}$.

The object here is only to show that a variety of topologies can be induced in $~\mathbf{X}$.

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Finally an application to thermodynamics is presented, there the concept of homogeneity is a basic one to define measures in the thermostatic space.

Entropy is assimilated to a \mathcal{C} -PROX.

Future developments are being pursued in, fundamentally, three directions: Other kinds of \mathcal{H} and \mathcal{G} , eventually non-fuzzy, new shapes for Γ i.e., plury-modal characterizers and new applications.

I - Symbology X is a set. x , y , z , ω elements of X . Λ is a lattice. ξ , δ , Σ , $\Gamma(\alpha)$ are elements of Ω . R non-negative real numbers. lpha , eta , eta , eta , eta , eta are elements of IR* $Q \in \mathcal{N}^{R^*}$ $\mathcal{R}: X \times X \rightarrow \mathcal{G} \quad (\infty, y) \mapsto \Gamma_{x, y}$ suprėmum of lattice 1 0 infimum of lattice V suprimum (or maximum), is a connective in $\mathcal N$; sup (max) Λ infimum (or minimum), is a connective in \mathfrak{N} ; inf (min). چ order on ${\cal N}$. equivalent to \succ and not \equiv . ۲ Ax AXIOM DEF DEFINITION REMARK REM ΤH THEOREM

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II - The Axioms

In the sequel X is a set and ${\mathcal N}$ is a complete, distributive lattice, and, in certain cases, an involutive operator (;) can be defined such that the order of the lattice is inverted; (\leq , \wedge , \vee , $_$ and $\stackrel{1}{2}$) are the order, the connectives inf (Λ) and sup (\vee), the infimum and the supremum of the lattice respectively, ${\cal G}$ is a subset of the family of the functions defined over $R^{\dagger} \equiv [O_{, \dagger} \infty)$ with values in \mathcal{N} , i.e. $\mathcal{G} \subseteq \mathcal{N}^{R^{\dagger}}$. At last it is given a function $\mathcal{R}: X \times X \rightarrow \mathcal{G}_{,}(x, y) \longrightarrow \Gamma_{xy}$. DEF.po $\mathcal{N} = \mathcal{N} \setminus \{ 0 \}$ DEF:pl A ternary (x, n, n) is said a \mathcal{C} -proximity, if the following axioms are satisfied: $\int \frac{\Omega }{L_{\text{Sublimber}}} \int \frac{\Delta L_{\text{Sublimber}}}{A \times 1a: \text{ If }} \int_{xy} \frac{(a)}{\epsilon} \left(Q \right), \text{ then } \int_{xy} \left(0 \right) = 0$ DEF. p2: $\Gamma_{xy} \triangleq V [\Gamma_{xy} (\alpha), \forall \alpha \in \mathbb{R}^+]$ DEF. p3: $\{ \alpha^* \} \triangleq \{ \alpha \in IR^* | \Gamma_{xy}(\alpha) = \Gamma_{xy}^* \}$ 1 -----Ax.1b: $\{\alpha^*\}$ is a closed interval $[(\gamma^*_{xy}, \gamma^*_{xy}])$, eventually a singleton, where $(\gamma^*_{xy}, \gamma^*_{xy}) \in IR$ and $(\gamma^*_{xy} \leq \gamma^*_{xy})$ Ax.lc: If $f_{xy} \in Q$, then $f_{xy}(\mathcal{A}_{\alpha}) \preceq \Gamma(\mathcal{A}_{k})$ whenever O & da & de < 4x . and Pay (dc) > Pay (dd) whenever Yay Ldc Edd , Xa, X&, Xc, Xd EIR* Ax. 1d: If $\Gamma_{xy} \in \mathcal{G}$, $\lim_{d \to \infty} \overline{\Gamma}(d) \rightarrow 0$ i.e. $\forall \varepsilon \in \Omega^{+}, d_{0} \in \mathbb{R}^{+}, \forall d \ge d_{0}, \Gamma(d) \prec \varepsilon$ Ax.le: If $\int_{xy} \in Q$, $\int_{xy} \langle A \rangle$ has, at most, a countable number of discontinuities. REM/O The concept of point of discontinuity for monotone functions is the usual one in lattice theory.

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	$\underline{\text{DEF} \cdot p4}: \mathcal{D}_{\Gamma_{xy}} \equiv \left\{ \mathcal{A} \mid \Gamma_{xy} \text{is discontinuous at } \mathcal{A} \right\}$
× ×/ /	<u>REM.1</u> : $\begin{bmatrix} \psi'_{xy}, \psi'_{xy} \end{bmatrix}$ will sometimes be represented by $A \left(\int_{xy} j = \int_{xy}^{x} \right)$ see REM.6
	<u>REM.2</u> : \int_{xy} is the characterizer or membership function of the fuzzy set just defined, and the shape of \int_{xy} suggests the name of "unimodal characterizer".
	on A
	<u>Ax.2a</u> : \mathcal{R} is symmetric, i.e. $\forall x, y \in X$, $\mathcal{R}(x, y) = \mathcal{R}(y, x) \Leftrightarrow \int_{xy} (d) = \int_{yx} (d) , \forall d$
l'td.	<u>Ax.2b</u> : $\forall x \in X$, $\Re(x, x) \iff \int_{x, x} (a) = 0 \in \mathcal{N}$, \angle
]. [let	$\frac{\text{DEF.p5: Let } x, y, z \in X, d, v \in [R^{+}, d \in [0, v] \subseteq IR^{+}]}{\text{define } \int_{xyz} : IR^{+} \longrightarrow \mathcal{N}, v \mapsto \int_{xyz} (v),$ and $\int_{xyz} (v) \triangleq V_{a} \cdot \Lambda [\int_{xy} (d), \int_{yz} (v - d)]$
	REM.3: This operation is represented by 🖨 , i.e.
	lxyz = lxy Olyz
	<u>DEF. D6</u> : $\int_{xyz}^{x} \neq V_{\sigma} \left[\int_{xyz}^{z} (\sigma) \right] \forall \sigma \in \mathbb{R}^{+} \in \mathcal{N}$
	$\underline{\text{DEF}} : \{ \mathcal{T}^{*} \} \triangleq \{ \mathcal{T} \in \mathcal{N} \mid \mathcal{I}_{xyz}(\mathcal{T}) = \mathcal{I}_{xyz}^{*} \}$
	<u>Ax.2c</u> : $\{\gamma^{x}\}$ is a closed interval $[\gamma^{x}_{xyz}, \gamma^{x}_{xyz}]$, eventually a singleton, where $\gamma^{x}_{xyz}, \gamma^{x}_{xyz} \in IR^{+}$ and $\gamma^{x}_{xyz} \leq \gamma^{x}_{xyz}$
XXXX	Ax.2d: For x, ZEA and VYEA, it is × Txz & Txyz, Yxz & Yxyz and Yxz & Yxyz
Jif Xand Qare topologia	Ax.2e: \mathcal{R} is a continuous function $\underline{/}$
spaces,	REM.4: Ax.2d is viewed as fuzzy triangular law and if "crisp" sets and e "crisp" triangular law are used then an "écart" is obtained.
	REM.5: In some applications, Ax.2b or Ax.2e or both are deleted.

III - <u>Theorems</u>

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$$\begin{array}{l} \begin{array}{l} \begin{array}{l} \underline{\mathsf{DEF}} & \underline{\mathsf{A}} \in [\Gamma_{\mathbf{x}\mathbf{y}} \times \mathcal{E}] \triangleq \left\{ d \in [\mathsf{R}^{+} \mid \Gamma_{\mathbf{x}\mathbf{y}}(d) \not\approx \mathcal{E} \mid \mathcal{E} \in \Omega \right\} \\ \text{ is the "crisp" set of } \Gamma_{\mathbf{x}\mathbf{y}}^{-} \ \text{ indexed by } \mathcal{E} \quad \underline{\mathsf{A}} \ \text{ is a closed} \\ \\ \begin{array}{l} \text{ interval on } [\mathsf{R}^{+} \\ & \underline{\mathsf{A}} \left(\Gamma_{\mathbf{x}\mathbf{y}}_{\mathbf{y}} \not\approx \mathcal{E} \right) = \begin{bmatrix} \Psi_{\mathbf{x}\mathbf{y} \not\approx \mathcal{E}} \ \mathbf{y} \quad \Psi_{\mathbf{x}\mathbf{y} \not\approx \mathcal{E}} \end{bmatrix} \leq [\mathsf{R}^{+} \\ \end{array} \\ \begin{array}{l} \begin{array}{l} \underline{\mathsf{OEF}} & \underline{\mathsf{P}}^{\mathbf{y}} \\ \underline{\mathsf{A}} \in [\mathsf{R}^{+} \mid \Gamma_{\mathbf{x}\mathbf{y}}(d) \not\approx \mathcal{E} \ \mathcal{E} \in \Omega \end{array} \right\} \\ & \underline{\mathsf{a}} \left\{ d \in [\mathsf{R}^{+} \mid \Gamma_{\mathbf{x}\mathbf{y}}(d) \not\approx \mathcal{E} \ \mathcal{E} \in \Omega \end{array} \right\} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{a}} \left\{ d \in [\mathsf{R}^{+} \mid \Gamma_{\mathbf{x}\mathbf{y}}(d) \right\} & \mathcal{E} \ \mathcal{E} \in \Omega \end{array} \right\} \\ \\ \begin{array}{l} \underline{\mathsf{where}} \end{array} \\ & \underline{\mathsf{stands}} \ \text{for } (\not\Rightarrow \ \mathsf{and not} \ \equiv 1. \ \mathsf{A} \ \text{ is the "crisp" set of} \\ \hline \Gamma_{\mathbf{x}\mathbf{y}} \ \mathcal{I} \ \text{ indexed by } \mathcal{E} \ \text{ and on open interval on } [\mathsf{R}^{+} \ \mathcal{R} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{REN.5}} \ \text{ referring to } \ \mathsf{REN.1}, \ \left\{ d^{+} \right\} \ \text{ can be represented by } \ \mathtt{A} \ \left(\Gamma_{\mathbf{x}\mathbf{y},\mathbf{y} = \Gamma_{\mathbf{x}\mathbf{y}} \right) \\ \end{array} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{REN.5}:} \ \text{ referring to } \ \mathsf{REN.1}, \ \left\{ d^{+} \right\} \ \text{ can be represented by } \ \mathtt{A} \ \left(\Gamma_{\mathbf{x}\mathbf{y},\mathbf{y} = \Gamma_{\mathbf{x}\mathbf{y}} \right) \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{REN.5}:} \ \text{ referring to } \ \mathsf{REN.1}, \ \left\{ d^{+} \right\} \ \mathsf{ can be represented by } \ \mathtt{A} \ \left(\Gamma_{\mathbf{x}\mathbf{y},\mathbf{y} = \Gamma_{\mathbf{x}\mathbf{y}} \right) \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{REN.5}: \ \text{ referring to } \ \mathsf{REN.1}, \ \left\{ d^{+} \right\} \ \mathsf{ can be represented by } \ \mathtt{A} \ \left(\Gamma_{\mathbf{x}\mathbf{y},\mathbf{y} = \Gamma_{\mathbf{x}\mathbf{y}} \right) \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \end{array} \\ \begin{array}{l} \underline{\mathsf{REN.5}: \ \text{ referring to } \ \mathsf{REN.1}, \ \left\{ d^{+} \ \mathbf{L} \ \mathbf$$

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$$\int_{xyz}^{x} (\sigma) \not\leq \varepsilon^{\circ} \text{ and } \int_{xyz}^{x} \not\equiv V \int_{xyz}^{z} (\sigma)$$

results $\int_{xyz}^{x} = \varepsilon^{\circ}$

PROOF Ax.1c: If
$$\forall_{i} \leq \forall_{z} \leq \forall_{xyz}^{*} = \mathcal{E}^{\circ}, \forall \geq d_{z}$$
 and $\forall \geq d_{1}$
it is $\int_{xy}^{1} (d_{2}) \And \int_{xy}^{1} (d_{1}), Ax.1c$,
and $\int_{yz}^{1} (\forall - d_{1}) \succcurlyeq \int_{yz}^{1} (\forall - d_{2}), Ax.1c$,
then, $\forall \land [\int_{xy}^{1} (d), \int_{yz}^{1} (\forall_{z} - d)] \succcurlyeq$
 $\Rightarrow \forall \land [\int_{xy}^{1} (d), \int_{yz}^{1} (\forall_{1} - d)],$
and $\int_{xyz}^{1} (\forall_{2}) \succcurlyeq \int_{xyz}^{1} (\forall_{1})$

If
$$\Psi_{xyz} = \mathcal{E} \leq \mathcal{E}_3 \leq \mathcal{E}_4$$

then $\Gamma_{xyz} (\mathcal{E}_3) \not\models \Gamma_{xyz} (\mathcal{E}_4)$

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$$\begin{array}{l} \underline{\mathsf{PRODF} \ \mathsf{Ax.1d}: \ \mathrm{If} \ \mathfrak{F} \to \infty, \ \mathfrak{A} \ \ \mathrm{and/or} \ (\mathfrak{F} - \mathfrak{A}) \to \infty \\ & \text{and} \ \Lambda \left[\left[\Gamma_{xy} \left(\mathfrak{A} \right), \ \Gamma_{yz} \left(\mathfrak{F} - \mathfrak{A} \right) \right] \to \mathcal{O} \in \mathcal{M} \\ & \text{or} \ \left[\Gamma_{xy} \left(\mathfrak{A} \to \infty \right) \to \mathcal{O} \in \mathcal{M} \\ & \text{and} \ \left[\Gamma_{yz} \left(\left(\mathfrak{F} - \mathfrak{A} \right) \to \infty \right) \right] \to \mathcal{O} \in \mathcal{M} \\ & \text{and} \ \left[\Gamma_{xyz} \in \mathcal{N}^{\mathsf{IR}^{\dagger}}, \ \mathrm{by} \ \mathrm{construction}. \\ \\ \underline{\mathsf{REM} \ \mathfrak{f}: \ \Gamma_{xyz} \in \mathcal{N}^{\mathsf{IR}^{\dagger}}, \ \mathrm{by} \ \mathrm{construction}. \\ \\ \underline{\mathsf{REM} \ \mathfrak{f}: \ \Gamma_{xyz} \left(\mathfrak{F} \right) \neq \mathcal{O} \in \mathcal{N}, \ \mathrm{in} \ \mathrm{general} \ \mathrm{and} \ \mathsf{Ax.2d} \ \mathrm{states} \ \mathrm{that} \ \Gamma_{xzz}^{\star} \succcurlyeq \Gamma_{xyz}^{\star} \\ & \mathrm{but} \ \mathrm{the} \ \mathrm{set} \ \mathfrak{F}^{\star} \ (\mathrm{see} \ \mathrm{DEF.p7}) \ \mathrm{may} \ \mathrm{not} \ \mathrm{contain} \ \mathcal{O} \ (\mathrm{zero}). \\ \\ \\ \underline{\mathsf{PRODF} \ \mathsf{Ax.1e}: \ \mathrm{As} \ \mathrm{both} \ \Gamma_{xy} \ \mathrm{and} \ \Gamma_{yz} \ \ \mathrm{have} \ \mathrm{at} \ \mathrm{most} \ \mathrm{countable} \ \mathrm{discontinuities}, \\ \hline{\Gamma_{xyz}} \ \mathrm{has} \ \mathrm{too} \ \mathrm{at} \ \mathrm{most} \ \mathrm{countable} \ \mathrm{discontinuities}. \end{array}$$

 $\begin{array}{ccc} \underline{19} & \underline{\text{REM 6}}: & \hline{\mathbf{1}}_{\mathbf{X}\mathbf{Y}\mathbf{Z}} & \text{satisfies Ax.la to Ax.le, and the operator } & \text{is closed for} \\ & \text{countable applications and if } & \hline{\mathbf{1}}_{\mathbf{X}\mathbf{Y}} & \text{and } & \hline{\mathbf{1}}_{\mathbf{Y}\mathbf{Z}} & , & \forall \mathbf{x}, \mathbf{y}, \mathbf{z} \in \mathbf{X} & , \\ & \underline{1}_{\mathbf{1}}^{\gamma} & \\ & \text{have no discontinuities for arbitrary applications.} \end{array}$

IV - Remarks on Topology

As referred in the introduction, the main goal here is to point out the adaptability of the $\,\mathcal{C}$ -prox..

The topology endowed in X depends up on each application. A natural suggestion is that the topology in X be such that the mapping $\mathcal{R}: X \times X \rightarrow \mathcal{N}^{\mathsf{R}^+}$ be continuous, when $\mathcal{N}^{\mathsf{R}^+}$ is endowed with a topology adjusted to the application.

It is remembered that the concept of discontinuity of f_{xy} supposes already that some kind of topology is in \mathcal{N} (see REM.8). Any finer topology than the initial one is suitable. Work is being done now on the concept of fuzzy balls to endow in X with a suitable topology.

Many authors have suggested topologies for \mathcal{N} or $\mathcal{N}^{\mathbb{R}^{+}}$, e.g. See Ref: 1, $\mathcal{I}_{\mathcal{I}}$ 3, 5, 7, 8, 9, 10, 11, 18. V - APPLICATION OF PROXIMITIES TO THERMODYNAMICS

(5)
$$\frac{\mu_{i}(\tau_{\alpha})}{\mu_{i}(\tau)} = \frac{\mu_{i}(\tau_{\alpha})}{\mu_{i}(\tau)}, \quad \forall \tau_{\alpha} \in \pi_{\overline{v}}(\tau)$$
$$\forall \mu_{i}, \mu_{i} \in \varphi$$

DEF. 2 The fineness γ of a partition $\pi_{\sigma}(\tau)$ is defined by (6):

(6)
$$\max\left[\frac{\mu_{i}(T_{\alpha})}{\mu_{i}(T)}, \forall T_{\alpha} \in TT_{\sigma}(T)\right] = p_{TT_{\sigma}}(T)$$

Taking in consideration (4) and (5), $\gamma_{\tau_{\tau}}(\tau)$ is independent of μ_{ι} .

$$\begin{array}{l} \underline{\text{DEF. 3}}: \ensuremath{ \baselineskip}{T_{T}} & \equiv \left\{ \ensuremath{ \baselineskip}{T_{T}}(\tau) : \ensuremath{ \baselineskip}{T_{T}}(\tau) \leq \ensuremath{ \baselineskip}{P_{T}} \right\} & \dots & (\mp) \end{array}$$
is the set of all partitions of T that have finenesses less than $\ensuremath{ \baselineskip}{P_{T}} \cdot \ensuremath{ \baselineskip}{T_{T}} \ensuremath{ \baselineskip}{I_{T}} \left\{ \ensuremath{ \baselineskip}{T_{T}} & \ensuremath{ \baselineskip}{P_{T}} + \ensuremath{ \baselineskip}{T_{T}} \ensuremath{ \baselineskip}{I_{T}} \ensuremath{ \baselineskip}{P_{T}} \ensuremath{ \baselineskip}{I_{T}} \ensuremath{ \baselineskip}$

for \forall , \top , \top , $\overleftarrow{} \in \mathcal{G}$ and $\forall \forall$ -homogeneous. Then \forall is an "adequate" set of linearly independent real measures (\neg -measures) for $\overrightarrow{\mathcal{G}}$.

Note that any other real measures on $\,\mathcal{T}\,$ can be expressed as a linear homogeneous function of degree 1 of the measures belonging to $\,\mathcal{Y}\,$.

B) <u>An Axiomatic for Thermostatic</u>

Ax. 1: All thermodynamic system is an universal class of sets \mathcal{G}^{*} , \mathcal{P}^{*} -homogeneous and \mathcal{P} is an "adequate" set of linearly independant real measures (∇ -measures) for \mathcal{G}^{*} . Card $\mathcal{P} = \mathcal{N}$, finite. For $\mathcal{V} \longrightarrow \mathcal{O}$ (zero), $\forall \mu_{i} \in \mathcal{P}$, μ_{i} (\mathcal{T}) is continuous on \mathcal{G}^{*} . Ax. 2: There are two real measures (∇ -measures), entropy μ_{o} and internal energy μ_{μ} . If $\mathcal{P} \equiv \{\mu_{o}, \mu_{i}...\mu_{i}\}$ then $\mu_{\mu} = \mathcal{F}[\mathcal{P}]$ and $\mathcal{P} \cup \mu_{\mu}$ is a N + 1 Euclidean Convex Space. μ_{μ} and μ_{o} are dual functions, exgratis: $(\min \mu_{\mu})\mu_{o} = \operatorname{const.} = (\max \mu_{o}) \mu_{\mu} = \operatorname{const.}$

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Note 1 : Continuous trajectories (lines) can be described on the surface

$$\mu_{\mathcal{M}} - \mathbb{F}\left[\mu_{s}, \mu_{i} \dots, \mu_{s}\right] = 0 \text{ if } P \longrightarrow 0.$$

Note 2 : The thermodynamic space is not metrisable, but a proximity can be defined, as it will be shown in C).

C)Proximity in thermodynamics

1) Reversible and irreversible trajectories**

In all trajectories (reversible or otherwise) the starting state τ^* and finishing state τ^* belong to $\mathcal{G}^* \not_{\gamma} \gamma$ -homogeneous.

If all the other intermediate states belong to \mathcal{Z}' then the trajectory is declared reversible, if not irreversible.

A general irreversible trajectory is symbolised in the following fashion:

 T^{∞} , T^{y} , T^{z} , T^{y} being respectively the starting and finishing point and T^{∞} , $T^{y} \in G^{\times}$.

Reversible trajectories are represented as follows: $\neg \xrightarrow{\sim} \neg \xrightarrow{\vee}$

2) Definition of a Proximity on \overline{C}^{\star}

 $\mathcal{H} = \mathcal{H} \mu_{i}, \forall \mu_{i} \in \Psi$ (Cartesean Product) The proximity of two states $T^{\infty} = T^{9}$ is given by:

> $\mathcal{R}(x,y) = \Gamma x y (\alpha)$ where $\Gamma x, y (\alpha) \in \mathcal{G}^{T} \subseteq \mathcal{G}$ (defined in 2 e)

 ${\boldsymbol{G}}^{\intercal}$ satisfies to the following conditions:

a) $C^{T} \in C$ b) $\Gamma'(o) = 0$

c) The non-decreasing branch starts at \propto = 0. Thus the general aspect of Γ is the following:



 \neq corresponds to the region of the "real" trajectories, the most plausible, and the reversible trajectories to $\measuredangle = 0$, $\Gamma(o) = 0$, which can be interpreted as "impractible" because $\Gamma(o) = 0$.

Comments:

- 1) All reversible trajectories are equiproximate $\ll = 0$, and their likelihood, $\int (\alpha = 0) = 0$, is zero, physically ideal.
- 2) All irreversible trajectories correspond to $\ll > \circ$ and their likelihood is positive $\vec{\Gamma}(\alpha > \circ) > \circ$.
- 3) The most likely proximity corresponds to the zone where $\int (lpha)$ is maximum.
- 4) Considering the two trajectories, $T^{x} \swarrow T^{y} \swarrow T^{z} \Longrightarrow \Gamma_{xyz}(\sigma)$ and $T^{x} \swarrow T^{z} \iff \Gamma_{xz}(\sigma)$. The zone of higher likelihood is shifted to greater values of α in Γ_{xyz} than in Γ_{xz} .

D) Entropy Production

If α is interpreted as entropy production, $\alpha \equiv \delta S$, then $\int_{xy}^{\infty} (\alpha) \equiv \int_{xy}^{\infty} (\delta S)$ and some simple conclusions can be taken:

- In a reversible trajectory (process), $\varkappa = 0$ then $\int_{\infty y}^{\infty} (o) = 0$, the likelihood of such a process is nill,
- The max $\left[\int_{xy} (\alpha) \right] = \int^{x}$ corresponds to the entropy production $55 = \alpha$ more likely to occur.
- The set $\left(\alpha : \Gamma_{xy}(\alpha) \geqslant \delta \leqslant \Gamma^* \right) \equiv \left[\delta_{\alpha}, \delta_{k} \right]$ is an interval of occurence of trajectories which are δ -likely to occur.

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and

 $\Gamma_{xz} \ge \Gamma_{xyz}$ $S_{a}^{*} \le \nabla_{a}^{*}$ and $S_{b}^{*} \le \nabla_{b}^{*}$

which means the entropy production in the process $T^{*} \longrightarrow T^{2}$ is greater than in the process $T^{*} \longrightarrow T^{2}$, for the same likelihood (or level of membership),

- If, max $\left[\int_{xy_2} (v) \right] = \int_{xy_2}^{x}$ and max $\left[\int_{x^2} (v) \right] = \int_{x^2}^{x}$

 $\begin{bmatrix} \nabla a & , \nabla b \end{bmatrix} \equiv \left\{ \nabla : \begin{bmatrix} x & yz & (\nabla) \\ x & yz \end{bmatrix} \equiv \left\{ \sigma : \begin{bmatrix} x & yz \\ x & z \end{bmatrix} \right\}$

E) Time and Entropy Production

If time t is considered a monotonous function of 1/2 = 1/25, some interesting interpretations are possible.

- a) If $\alpha = 0$ then $\pm = \infty$. A process that would take ∞ time for completion would be eventually reversible.
- b) The typical t^* (or the most likely time) would correspond to $\Gamma^*(\max \Gamma(\alpha))$.
- c) As $\prec \rightarrow \infty$, $t \rightarrow o$, and $\Gamma(\checkmark) \rightarrow o$. this could be interpreted as follows: when the time of the process is less than t^* , then the likelihood of the process would diminish tending to zero with $t \rightarrow o$.

Conclusion

Space $\mathfrak{X} \equiv \mathfrak{P}$ can be topologically structured with a class $\mathcal{C} \leq \mathcal{C}$ of proximities and some form of a <u>fuzzy distance</u>. Proximity, between thermodynamic states can be defined.

Entropy production is a monotonous function of \prec , eventually $\measuredangle \equiv SS$. Time is an inverse function of SS.

V - Conclusion

The flexibility of Proximities gives a certain amount of freedom in the structuring of X, by choosing a convenient $\mathcal{R}: X \longrightarrow B \in \mathcal{N}^{\mathbb{R}^{+}}$.

 R^+ can be substituted by a suitable compact, various $\mathcal N$ and respective topologies are presented in the literature and can be used.

In some applications to graphs and hyper-graphs, one can delect Ax.2b and when the real system projected in $\,\times\,$ is non-homogeneous, Ax.2c can be supressed.

Work is being done on proposing suitable topologies for ${\mathcal N}$ and ${\mathsf X}$.

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